

Significance calculation and a new analysis method in searching for new physics at the LHC

Yongsheng Gao^a, Liang Lu, Xinlei Wang

Southern Methodist University, Dallas, TX 75275-0175, USA

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Abstract. The LHC experiments have great potential in discovering many possible new particles up to the TeV scale. The significance calculation of an observation of a physics signal with known location and shape is no longer valid when either the location or the shape of the signal is unknown. We find the current LHC significance calculation of new physics is over-estimated and strongly depends on the specifics of the method and the situation it applies to. We describe general procedures for significance calculation and comparing different search schemes. A new method uses maximum likelihood fits with floating parameters and scans the parameter space for the best fit to the entire sample. We find that the new method is significantly more sensitive than current method and is insensitive to the exact location of the new physics signal we search.

1 Introduction

The Large Hadron Collider (LHC) at CERN will open a new frontier in particle physics due to its higher collision energy and luminosity as compared to the existing accelerators. The general-purpose ATLAS and CMS experiments at the LHC will employ precision tracking, calorimetry and muon measurements over a large solid angle to identify and measure electrons, muons, photons, jets and missing energy accurately. Therefore, they have great physics potential in discovering many possible new particles. Among them are the Standard Model (SM) Higgs boson, supersymmetric (SUSY) and other new particles beyond the SM. All of them can have masses in a very large range up to the TeV scale. The significance calculation in searching for and observation of a physics signal with known location and shape is no longer valid when either the location or the shape of the signal is unknown. This will be the case for many of the possible new physics signals at the LHC.

In Sect. 2, we give a short review of the significance calculation and current analysis strategy in High Energy Physics (HEP) and at the LHC. Using a signal with known shape but unknown location as an example, we discuss in detail in Sect. 3 the problems of the current significance calculation. We then describe general procedures for significance calculation and comparing different search schemes in Sect. 4. In Sect. 5, we describe a new analysis method and compare it with the current “Sliding-Window” approaches following these procedures. Detailed comparison results are also given in this Section. Summary and discussion are given in Sect. 6. In this note we limit ourselves to the significance calculation and analysis method used in searching for an individual decay mode of new physics signals.

2 Review of significance calculation and current analysis strategy to search for new physics at the LHC

In the field of HEP, a common strategy to detect a physics signal is to search for an excess of events in a certain region of a kinematic observable. The observation probability is given by Poisson statistics:

$$P(n, B) = \frac{e^{-B} B^n}{n!} \quad (1)$$

where B is the number of the expected events to be observed in the region, and n is the number of the observed events in this region. When B is large (over 25, for instance), the significance of an observation can be approximated well by S/\sqrt{B} of Gaussian statistics, where $S = n - B$.

In HEP, the significance of an observation is defined by the probability that such an observation is due to statistical fluctuation of background events. When we claim an observation has a significance of 5σ [1], the common criterion for a HEP discovery, the probability that the claimed discovery is due to statistical fluctuation of background events, known as the Type I error rate in statistics, needs to be less than 2.9×10^{-7} . The background fluctuation probabilities which define the 1σ to 5σ significances in HEP are shown in Table 1.

If the expected mass spectrum of a physics signal is a Gaussian distribution with standard deviation σ , the mass region used to calculate the observation significance of this signal is usually $\pm 2\sigma$ around the Gaussian mean. Including regions where the physics signal has little chance to show up only increases B and decreases S/\sqrt{B} . This is why the region in which to search for the signal and calculate significance is usually limited to $\pm 2\sigma$ around the Gaussian

^a e-mail: gao@mail.physics.smu.edu

Table 1. The definition of significance in HEP and the corresponding background statistical fluctuation probabilities

Significance	1σ	2σ	3σ	4σ	5σ
Probability that the observation of the excess of events is due to background statistical fluctuation	15.87%	2.28%	0.14%	3.2×10^{-5}	2.9×10^{-7}

mean, in order to maximize the discovery potential and observation significance. This approach has been widely and successfully used in many HEP experiments at CESR, Tevatron, LEP, KEK-B, PEP-II, etc. It is only valid when searching for and observation of a physics signal with known location and shape, i.e., when the kinematic region for the significance calculation is uniquely defined.

One of the new challenges for the ATLAS and CMS experiments is that we do not know the masses of the new particles we will be searching for. The current analysis method proposed for new particle searches at the LHC is to use a “Sliding-Window”, i.e., look for an excess of events in a series of narrow regions or windows over the entire available kinematic range. The location and width of each window is given by the expectations of the new particle with a specific mass and the corresponding width. The expected significances and discovery potential for new particle searches are only determined by the S and B values within these narrow windows [2–4].

3 Problem with the current LHC search method and significance calculation

There is a fundamental problem in the above significance calculation. The significance of an observation is defined according to the probability that such an observation is due to statistical fluctuation of background events, i.e., the Type I error rate. The current expected significance calculation is only correct if we know exactly the location and shape of the new physics signal we are searching for, and we use only one window to search and calculate the observation significance. In the “Sliding-Window” method, we search for an excess of events in any of the narrow windows over a wide kinematic range, but still use the S and B of each narrow window to calculate the significance of the observation. Therefore, the probability of observing a “significant” excess of events due to background statistical fluctuation in any window will be much higher [5]. This “false-positives” problem caused by multiple testing was recognized in statistics many years ago [6, 7].

We use simple simulations to demonstrate this problem. Assume that we search for a possible Gaussian signal with a standard deviation $\sigma=1.0$ but an unknown mean between 2.0 and 98.0, and that the expected distribution of the background is flat between 0.0 and 100.0. We generate 13,450,000 background-only Monte Carlo (MC) experiments (referred to as the “background-only sample”) with each experiment containing 500 events generated from a flat distribution between 0.0 and 100.0.

Table 2. The probability of observing at least one “1”, “2”, “3”, “4” and “5” σ -effect window in any background-only MC experiment using the “Sliding-Window” method with various step sizes

Significance (S/\sqrt{B})	“1” σ	“2” σ	“3” σ	“4” σ	“5” σ
Step Size = 16	70.89%	20.42%	1.522%	0.11%	0.002%
Step Size = 8	91.56%	35.25%	2.818%	0.20%	0.003%
Step Size = 4	99.72%	58.53%	5.380%	0.39%	0.007%
Step Size = 2	99.99%	77.86%	9.635%	0.73%	0.015%
Step Size = 1	100.0%	89.03%	14.86%	1.24%	0.027%
Step Size = 0.5	100.0%	94.33%	19.97%	1.83%	0.042%
Step Size = 0.2	100.0%	97.17%	25.42%	2.56%	0.064%
Step Size = 0.1	100.0%	98.01%	28.21%	2.98%	0.078%

We use a “Sliding-Window” with a fixed width of 4.0 and move the center of this fixed-width window from 2.0 to 98.0 with various step sizes of 16.0, 8.0, 4.0, 2.0, 1.0, 0.5, 0.2 and 0.1, respectively to search for an excess of events in any of the windows. The fixed width of 4.0 of the “Sliding-Window” corresponds to $\pm 2\sigma$ around the unknown Gaussian mean. The significance in any one window of a MC experiment is calculated by S/\sqrt{B} according to the current significance calculation, where n is the number of events in that window of the experiment, $S = n - B$, and $B = 20$.

The probabilities that we observe at least one window with $S/\sqrt{B} > 1, 2, 3, 4, 5$ (i.e. “1”, “2”, “3”, “4” and “5” σ according to the current significance calculation) in any of the background-only sample experiment are shown in Table 2. This probability is defined as the number of background-only MC experiments which contain at least one “1”, “2”, “3”, “4” and “5” σ -effect window divided by the total number of background-only MC experiments. From Table 2, we can see that the probabilities of positive observations are much higher than the Table 1 background fluctuation probabilities that define the significances in HEP. Furthermore, the probability of finding a signal of given significance increases as the step size of the “Sliding-Window” decreases, i.e., as more windows are scanned over the same kinematic range. While each individual window follows Poisson or Gaussian statistics reasonably well, the probability of observing an excess in any of the multiple windows is much higher than that for an individual window. Table 2 clearly shows the problems of the significance calculation in searching for new physics signals with an unknown location. It is due to the fact that we search for an excess of events over multiple narrow windows, but the

significance is still calculated according to an individual narrow window.

4 Procedures for significance calculation and comparing different analysis approaches

Each analysis approach or search scheme can be described by two measures. The Type I error rate measures how often false signals are claimed when there are only background events. The significance of an observation as defined according to this error rate is shown in Table 1. The Power or Sensitivity measures how often real signals can be found correctly when they are present. There is a correlation between the two measures, the Type I error rate increases with increasing sensitivity. Therefore, we need to set one of these two measures to the same value for different search schemes and compare the other measure, in order to quantitatively compare these search schemes.

We can see from Table 2 that the “significance” calculated by S/\sqrt{B} of the “Sliding-Window” method is highly over-estimated compared to the HEP significance definition. Furthermore, it strongly depends on the specifics of the search scheme and the situation it applies to, i.e. step size of the “Sliding-Window” used to scan the kinematic range, the total range of the kinematic region, etc. We need to evaluate the significance reported by each scheme so it truly reflects our significance definition. The procedures to calculate significance and compare different search schemes are as follows:

1. Use background-only MC experiments to evaluate the significance of all search schemes. After the evaluation, all the search schemes should be normalized to have the same Type I error rates, which follow the HEP significance definition.
2. Use signal-embedded MC experiments to evaluate the sensitivity of the search schemes. The search scheme with the higher sensitivity is the better one.

These procedures are applied to compare a new analysis method with the current “Sliding-Window” approaches in the following section.

5 A new analysis method and a comparison with the “Sliding-Window” approaches

An alternative approach is to apply an unbinned maximum likelihood scan method with floating parameters to the entire sample and search for the best fit to the sample over the entire parameter space [5]. It is intended to minimize the sensitivity of the significance to local fluctuations. We follow the procedures described in Sect. 4 to compare the current “Sliding-Window” approaches with this new method for this example [8].

1. We search for a possible Gaussian signal ($\sigma=1.0$ with unknown mean between 2.0 and 98.0) on top of a flat background in the 13,450,000 background-only MC

experiments (“background-only sample”) using each search scheme. We then evaluate the significance of each scheme so that it follows the HEP significance definition for the background-only sample.

2. We generate signal-embedded MC experiments and perform the same search using each search scheme. We then calculate the sensitivities of finding the embedded signal for each search scheme based on the significances defined by the background-only sample.

For the “Sliding-Window” approach in Step 1, we make a table which defines the new cutoff values of S/\sqrt{B} which follow the HEP significance definition for the background-only sample. Similarly for the new approach, we find out the values of the Maximum Likelihood fit output which corresponds to 1, 2, 3, 4 and 5σ for the background-only sample according to the HEP significance definition.

5.1 Significance evaluation of “Sliding-Window” approaches

We use the background-only sample to evaluate the significance of the “Sliding-Window” approach. For each experiment, we use a “Sliding-Window” with fixed width of 4 and move the center of the window from 2.0 to 98.0 with step sizes of 16, 8, 4, 2, 1, 0.5, 0.2 and 0.1, respectively, to search for the window with the maximum S/\sqrt{B} . For each step size, we plot the maximum S/\sqrt{B} for all the background-only sample. We then find the corresponding cutoff values on the plot which follow the HEP significance definition. For example, the maximum S/\sqrt{B} from “Sliding-Window” approaches with step sizes of 16, 4, 1, and 0.1 for the background-only sample are shown in Fig. 1. In the “Sliding-Window” approach with step size of 0.1, we find that 15.87% of the experiments have at least one window with $S/\sqrt{B} > 3.35$, and 2.28% of the experiments have at least one window with $S/\sqrt{B} > 4.02$. According to our HEP significance definition in Table 1, the experiments which contain windows with $S/\sqrt{B} > 3.35$ are defined as 1σ for the “Sliding-Window” approach with step size of 0.1 in this case. Similarly, the experiments which contain window with $S/\sqrt{B} > 4.02$ are defined as 2σ . The new S/\sqrt{B} cutoff values which follow our HEP significances definition for the “Sliding-Window” approaches with various step sizes are given in Table 3. The cutoff values are not continuous, because $S = n - B$, $B = 20$, and both n and S are integers.

5.2 Significance evaluation of the new analysis method

We use the same background-only sample to evaluate the significance for the new approach. In this specific example we search for a Gaussian signal ($\sigma=1.0$ with unknown mean between 2.0 and 98.0) on top of a uniform background. The Likelihood is then calculated as:

$$L(Y|\mu) = \prod_{i=1}^n P(y_i|\mu) \quad (2)$$

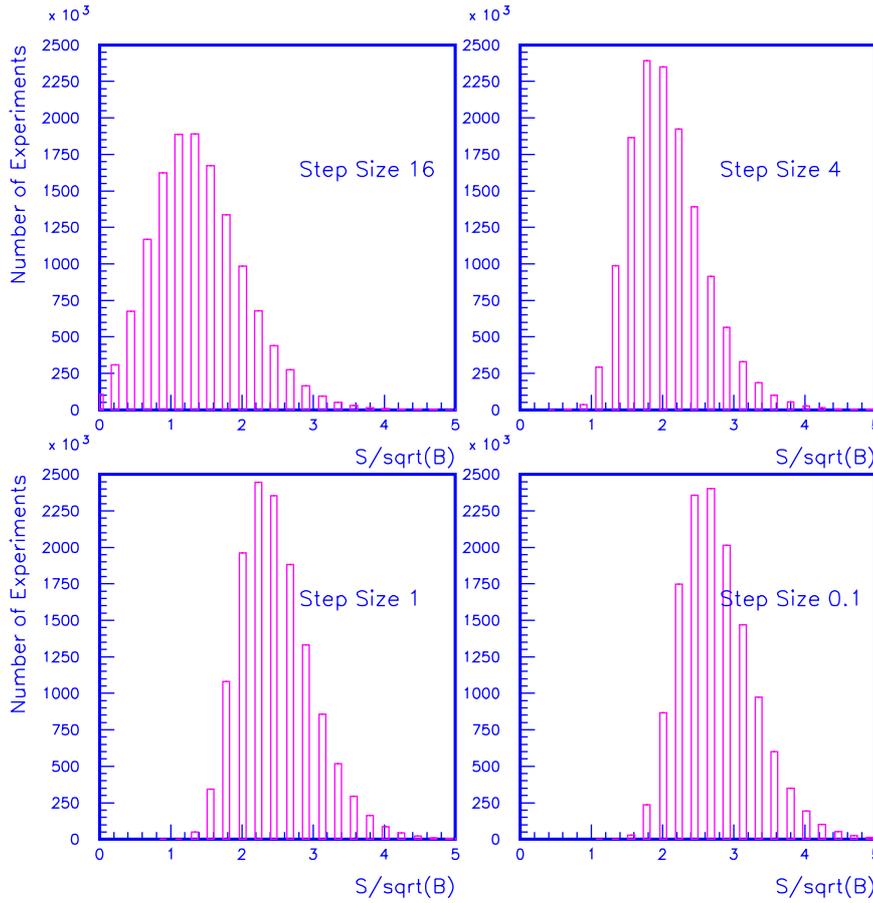


Fig. 1. The maximum S/\sqrt{B} from “Sliding-Window” approaches with step sizes of 16, 4, 1, and 0.1 for the 13,450,000 background-only MC experiments

Table 3. The S/\sqrt{B} cutoff values which correspond to the HEP significance definition for a “Sliding-Window” approach with different step sizes

Significance	1σ	2σ	3σ	4σ	5σ
Step Size = 16	2.01	2.90	3.80	4.91	6.03
Step Size = 8	2.23	3.13	4.02	5.14	6.26
Step Size = 4	2.68	3.35	4.24	5.36	6.48
Step Size = 2	2.90	3.57	4.47	5.36	6.48
Step Size = 1	2.90	3.80	4.47	5.59	6.48
Step Size = 0.5	3.13	3.80	4.69	5.81	6.93
Step Size = 0.2	3.13	4.02	4.91	5.81	6.93
Step Size = 0.1	3.35	4.02	4.91	5.81	6.93

where the Y are the data per experiment, and y_i is the individual data point in each experiment where $i = 1, 2, 3, \dots, n$ ($n = 500$). $P(y_i|\mu)$ is the normalized probability density of y_i as a function of the parameter μ which is the unknown mean of the Gaussian signal. The normalized probability density is given by:

$$P(y_i|\mu) = \frac{(1-p)}{100} + \frac{p}{\sqrt{2\pi}} e^{-\frac{1}{2}(y_i-\mu)^2} \quad (3)$$

where 100 is the normalization factor which guarantees that the integral of $P(y_i|\mu)$ over the range from 0.0 to 100.0 is equal to 1. p is the probability of the data point being

the Gaussian signal. Similarly, $(1-p)$ is the probability of the data point being the background. The optimization process attempts to find the μ parameter that maximizes $L(Y|\mu)$ for each experiment, or minimizes $-\log(L(Y|\mu))$. We use

$$\sum_{i=1}^{500} \log\left((1-p) + p \frac{100}{\sqrt{2\pi}} e^{-\frac{1}{2}(y_i-\mu)^2}\right) \quad (4)$$

as the maximum likelihood output to simplify the calculation.

For the MC experiment generation and maximum likelihood analysis we used the statistical computing software R [9, 10]. R is a language and environment for statistical computing and graphics. It is a GNU project developed at Bell Laboratories and provides a wide variety of statistical and graphical techniques (linear and nonlinear modeling, classical statistical tests, time-series analysis, classification, clustering, etc.) [11]. We have also tried other statistical software packages such as SAS [12], Matlab [13] and the HEP software package RooFit [14] to generate the MC experiments, and perform the maximum likelihood fits. The results with different analysis tools are all consistent. We decided to use R because it is faster than the other packages.

In order to find the best fit with a floating μ parameter for each experiment, we break down the μ parameter region from 2.0 to 98.0 into 96 equal intervals [15]. We perform one maximum likelihood fit for each interval to find the

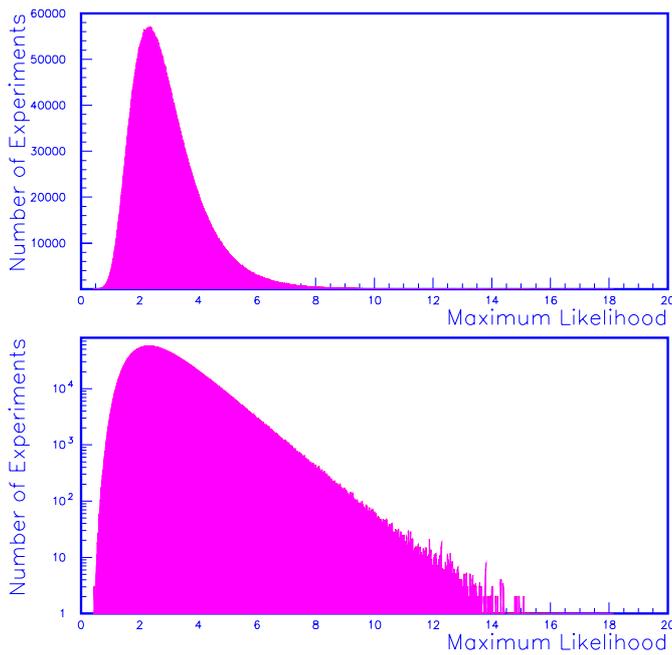


Fig. 2. The Maximum Likelihood output for 13,450,000 background-only MC experiments. The top plot is in linear scale while the bottom plot is in log scale

Table 4. The cutoff values of the maximum likelihood output for the new analysis method in this example

Fraction of background-only experiments below cutoff	Cutoff value	Significance
84.13%	4.00	1σ
97.72%	5.94	2σ
99.86%	8.71	3σ
99.9968%	12.48	4σ
$(1 - 2.9 \times 10^{-5})\%$	16.61	5σ

best fit. We then compare all 96 fits to find the overall best fit for the entire μ parameter space for this experiment. The maximum likelihood output of the best fit for the background-only sample is shown in Fig. 2. Because 84.13% of the background-only MC experiments have a maximum likelihood output below 4.00, the cutoff value for 1σ is set at 4.00. Similarly, 5.94 is set as the cutoff value for 2σ . The cutoff values for 1 to 5σ significances for the new analysis method in this example are given in Table 4.

After these evaluations, the significances reported by the “Sliding-Window” approaches and the new analysis method are all adjusted to follow the HEP significance definition for background-only sample. We can then compare the sensitivity of these approaches using the signal-embedded MC experiments.

5.3 Sensitivity comparison of “Sliding-Window” approaches with the new analysis method

We generate signal-embedded MC experiments to calculate the Power or Sensitivity of each approach. Each experiment contains a small number (5, 10, 15, 20, 25, 30,

and 35) of signal events generated according to a Gaussian distribution with $\sigma=1.0$ and a specific Gaussian mean (42.00, 46.00, 48.00, 49.00, 49.50, 49.75, 49.90, 49.95, and 50.00) [8]. Each signal-embedded experiment contains one set of these signal events embedded with 500 background events generated with a flat distribution between 0.0 and 100.0. A total of 630,000 signal-embedded MC experiments are generated, with 10,000 experiments for each set of Gaussian signal parameters. For example, 10,000 experiments each with 5 Gaussian signal events with Gaussian mean at 42.00 embedded into 500 background events are generated. The maximum likelihood output for each signal-embedded experiment is normalized to 500 events before comparing with the cutoff values from background-only experiments.

We use these signal-embedded experiments to calculate the sensitivity of the “Sliding-Window” approaches. We use a “Sliding-Window” of fixed width of 4.0 and move the center of this fixed-width window from 2.0 to 98.0 with various step sizes of 16.0, 8.0, 4.0, 2.0, 1.0, 0.5, 0.2 and 0.1, respectively, to search for the window with the maximum S/\sqrt{B} for each experiment. The success of finding the embedded signal in an experiment is defined as when the center of the most significant window is found within 1.0 of the Gaussian mean of the embedded signal events. The significance of this window is defined according to the cutoff values in Table 3.

We use the same signal-embedded experiments to calculate the sensitivity of the new analysis method. For each experiment, we break down the μ parameter region from 2.0 to 98.0 into 96 equal intervals. We perform one maximum likelihood fit for each interval to find the best fit for this interval and the corresponding μ value. We then compare all 96 fits to find the best overall fit for the entire μ parameter space and its corresponding μ value for each experiment. The success of finding the embedded signal in an experiment is defined as when the μ value of the best fit in the entire μ parameter space falls within 1.0 of the Gaussian mean of the embedded signal events. Similarly, the significance is defined according to the cutoff values in Table 4.

The sensitivity of the “Sliding-Window” approaches with various step sizes is very sensitive to the exact value of the Gaussian mean of the embedded signal events. Therefore, we choose the best-case and worst-case scenarios for each “Sliding-Window” approach with a specific step size. The best-case scenario corresponds to a case where the $\pm 2\sigma$ region of the embedded Gaussian signal falls exactly inside one of the “Sliding-Windows”. The worst-case scenario corresponds to a case when the embedded Gaussian signal falls exactly between two neighboring “Sliding-Windows”. In Table 5, we show the Gaussian means of the embedded Gaussian signals for the best-case and worst-case scenarios for various “Sliding-Window” step sizes.

For each set of the 10,000 signal-embedded MC experiments, we calculate how many times the embedded signals are correctly found by each approach with a significance greater than 1, 2, 3, 4 and 5σ according to the HEP significance definitions. For the “Sliding-Window” approach with a specific step size, two numbers are reported, according to the best-case and worst-case scenarios respectively. In

Table 5. The Gaussian means of the embedded Gaussian signal for the best-case and worse-case scenarios of the “Sliding-Window” approaches with various step sizes

Step Size	Gaussian Mean of embedded signal	
	Best-Case Scenario	Worst-Case Scenario
16	50.00	42.00
8	50.00	46.00
4	50.00	48.00
2	50.00	49.00
1	50.00	49.50
0.5	50.00	49.75
0.2	50.00	49.90
0.1	50.00	49.95

contrast, the new approach scans the parameter space and performs a maximum likelihood fit at each small interval to cover the entire parameter space to search for the best fit of the entire sample. Thus, it is not sensitive to the exact value of the Gaussian mean of the embedded signal. We find the number is independent of the exact location of the embedded Gaussian signal for the new analysis method.

The work and results for 5, 10, 15, 20, 25, 30, and 35 signal events embedded with Gaussian means at 42.00, 46.00,

48.00, 49.00, 49.50, 49.75, 49.90, and 50.00 are shown in [8]. The results for 10, 20, and 30 signal events embedded are shown in Tables 6, 7, and 8. We can see that the number of signal embedded experiments successfully found with a certain significance is much lower than what expected from S/\sqrt{B} calculations. This is a price we have to pay for not knowing the exact location of the signal. Furthermore, for the “Sliding-Window” method, the sensitivity strongly depends on the exact location of the embedded signal. If the step size is greater than 1, the embedded signals are totally missed for the worst-case scenarios. For step size of 1 or less, there are still significant differences in the sensitivities between the best-case and worse-case scenarios, depends on the step size of the “Sliding-Window” used to scan the kinematic range. In comparison, the new analysis method is independent of the exact location of the Gaussian mean of the embedded signal events. This is because the new method scans the entire parameter space for the best fit to the entire experiment. The maximum S/\sqrt{B} from “Sliding-Window” approaches with step sizes of 16, 4, 1, and 0.1 for the best-case scenario MC experiments each with 20 signal events embedded are shown in Fig. 3. Similarly, the maximum S/\sqrt{B} from “Sliding-Window” approaches with step sizes of 16, 4, 1, and 0.1 for the worst-case scenario MC experiments each with 20 signal events embedded are shown in Fig. 4. The maximum likelihood output of the best fits

Table 6. The number of signal embedded experiments successfully found with 1, 2, 3, 4 and 5σ significance in the 10,000 MC experiments each with 10 signal events embedded.

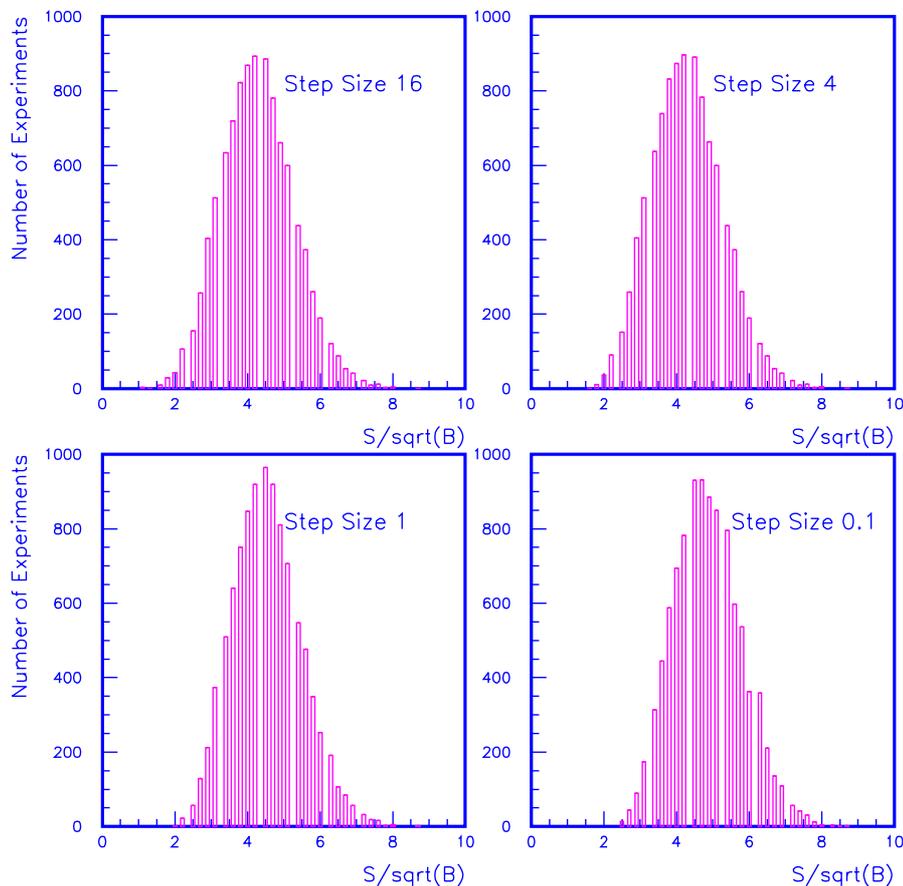
Significance Scenario	1σ	2σ	3σ	4σ	5σ
	Best/Worst	Best/Worst	Best/Worst	Best/Worst	Best/Worst
Step Size = 16	4795/0	1859/0	433/0	39/0	1/0
Step Size = 8	3897/0	947/0	271/0	23/0	1/0
Step Size = 4	1839/0	652/0	176/0	16/0	1/0
Step Size = 2	1691/0	619/0	107/0	16/0	1/0
Step Size = 1	1915/1031	537/274	143/68	9/1	1/0
Step Size = 0.5	2011/1728	807/706	156/120	11/2	2/0
Step Size = 0.2	2011/1728	562/484	97/70	11/2	2/0
Step Size = 0.1	1598/1465	599/548	106/78	13/2	1/0
New Approach	2328	819	153	17	2

Table 7. The number of signal embedded experiments successfully found with 1, 2, 3, 4 and 5σ significance in the 10,000 MC experiments each with 20 signal events embedded.

Significance Scenario	1σ	2σ	3σ	4σ	5σ
	Best/Worst	Best/Worst	Best/Worst	Best/Worst	Best/Worst
Step Size = 16	9834/0	8969/0	6302/0	2215/0	356/0
Step Size = 8	9694/0	7844/0	5435/0	1616/0	240/0
Step Size = 4	8932/0	7135/0	4542/0	1190/0	158/0
Step Size = 2	8439/0	6867/0	3642/0	1182/0	156/0
Step Size = 1	8086/4955	6365/3913	4032/2509	965/634	174/112
Step Size = 0.5	7795/7264	6812/6273	4020/3782	1046/965	130/100
Step Size = 0.2	7795/7264	6217/5747	3274/3061	1046/965	130/100
Step Size = 0.1	7538/7277	6314/6034	3416/3268	1119/1051	149/111
New Approach	9116	7406	4073	1209	206

Table 8. The number of signal embedded experiments successfully found with 1, 2, 3, 4 and 5σ significance in the 10,000 MC experiments each with 30 signal events embedded

Significance	1σ	2σ	3σ	4σ	5σ
Scenario	Best/Worst	Best/Worst	Best/Worst	Best/Worst	Best/Worst
Step Size = 16	10000/0	10000/0	9963/0	9215/0	5936/0
Step Size = 8	10000/0	9993/0	9917/0	8792/0	5079/0
Step Size = 4	9998/0	9987/0	9841/0	8249/0	4238/0
Step Size = 2	9904/0	9886/0	9633/0	8210/0	4232/0
Step Size = 1	9542/6589	9532/6584	9412/6490	7787/5444	4560/3285
Step Size = 0.5	9249/8784	9246/8781	9152/8665	7743/7331	3847/3683
Step Size = 0.2	9249/8784	9238/8778	9046/8558	7743/7331	3847/3683
Step Size = 0.1	9183/8945	9178/8934	9006/8742	7792/7603	4049/3944
New Approach	9985	9974	9723	8024	4332

**Fig. 3.** The maximum S/\sqrt{B} from “Sliding-Window” approaches with step sizes of 16, 4, 1, and 0.1 for the best-case scenario 10,000 MC experiments each with 20 signal events embedded

for MC experiments with 5, 10, 20 and 30 signal embedded are shown in Fig. 5. Compared to “Sliding-Window” approaches with a step size small enough not to miss the worst-case scenarios, the sensitivity of the new analysis method is significantly higher. This means that the new analysis approach is a significantly better and more sensitive scheme to search for new physics signals at the LHC than the current “Sliding-Window” method.

The analysis method described above performs a scan of the entire parameter space using unbinned maximum likelihood fits at every small interval of the parameter space. It is very CPU-intensive. The 13.45 million background-

only and 630,000 signal embedded MC experiments were generated and analyzed over several months with about 10 dual-CPU servers.

6 Summary and discussion

We have examined the significance calculation and analysis methods in searching for an individual decay mode of a new physics signal at the LHC. Unlike the search for a physics signal with known location and shape, the significance calculation for new physics signals with unknown

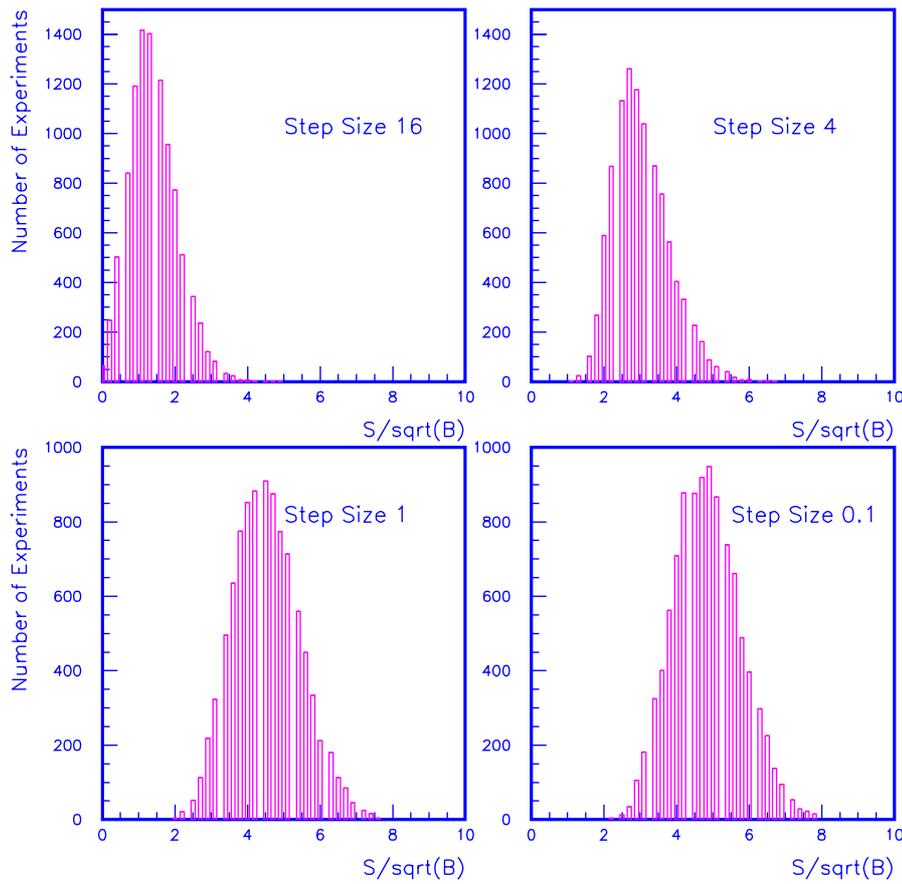


Fig. 4. The maximum S/\sqrt{B} from “Sliding-Window” approaches with step sizes of 16, 4, 1, and 0.1 for the worst-case scenario 10,000 MC experiments each with 20 signal events embedded

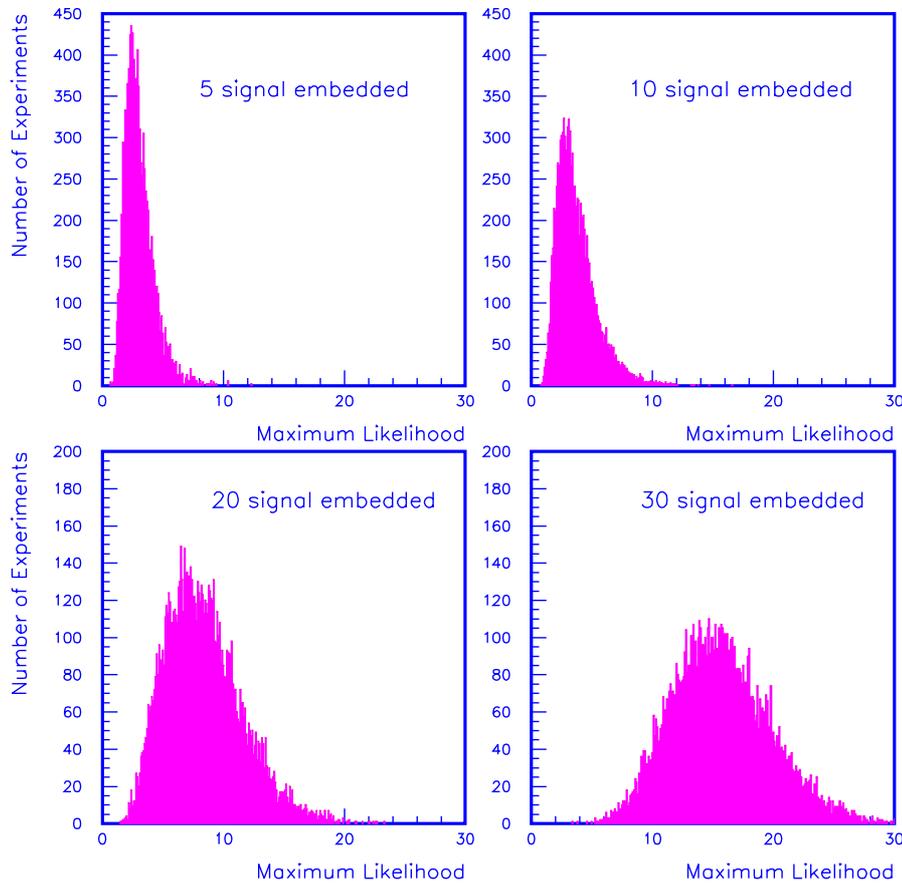


Fig. 5. The maximum likelihood output for MC experiments with 5, 10, 20, 30 signal events embedded respectively

location or shape strongly depends on the details of the search scheme and the situation it applies to. Using a signal with known shape but unknown location as an example, we have demonstrated that the significance calculation using the current “Sliding-Window” method at the LHC is over-estimated. This is because we search for an excess of events over multiple narrow windows, but the significance is still calculated according to an individual narrow window. The significance and sensitivity of the “Sliding-Window” method strongly depends on the specifics of the method and the situation it applies to, e.g. the step size of the “Sliding-Window” used to scan the available kinematic range, the total available kinematic range to search for the new physics signal, and the exact location of the new physics signal, etc.

We describe general procedures for significance calculation and comparing different search schemes. We have applied the procedures and compared the current “Sliding-Window” approaches with a new analysis method. The proposed new analysis method uses maximum likelihood fits with floating parameters and scans the parameter space for the best fit to the entire sample. We find the results of the new analysis method is independent of the location of the new physics signal and significantly more sensitive in searching for new physics signal than the current “Sliding-Window” approaches.

While the LHC experiments have great potential in discovering many possible new physics signals, we need to be extremely careful in evaluating the significance of an observation from the real LHC data. Because possible new physics can show up in many kinematic observables, over a very large kinematic range, the fluctuation probability of background events will be much higher. For individual decay modes of new physics signals, the expected significances in observing the new physics signal will be much smaller than current expectations [2–4]. Combining independent decay modes of the same new physics signal will be essential to establish the discovery of the new physics signal. Significant observations of the same new particle in independent decay modes at consistent locations will be the most effective way to establish the discovery of this new particle. Careful evaluation of the observation significance in each individual decay mode following the general

procedures described in this paper is the starting point, before we can evaluate the significance of the observations of independent decay modes.

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